

Efficient Estimation of Climate State and Its Uncertainty Using Kalman Filtering and Application to Policy Thresholds

J. Matthew Nicklas,^a Baylor Fox-Kemper,^a Charles Lawrence.^a

^a *Brown University, Providence, Rhode Island.*

Corresponding author: J. Matthew Nicklas, john_nicklas@brown.edu

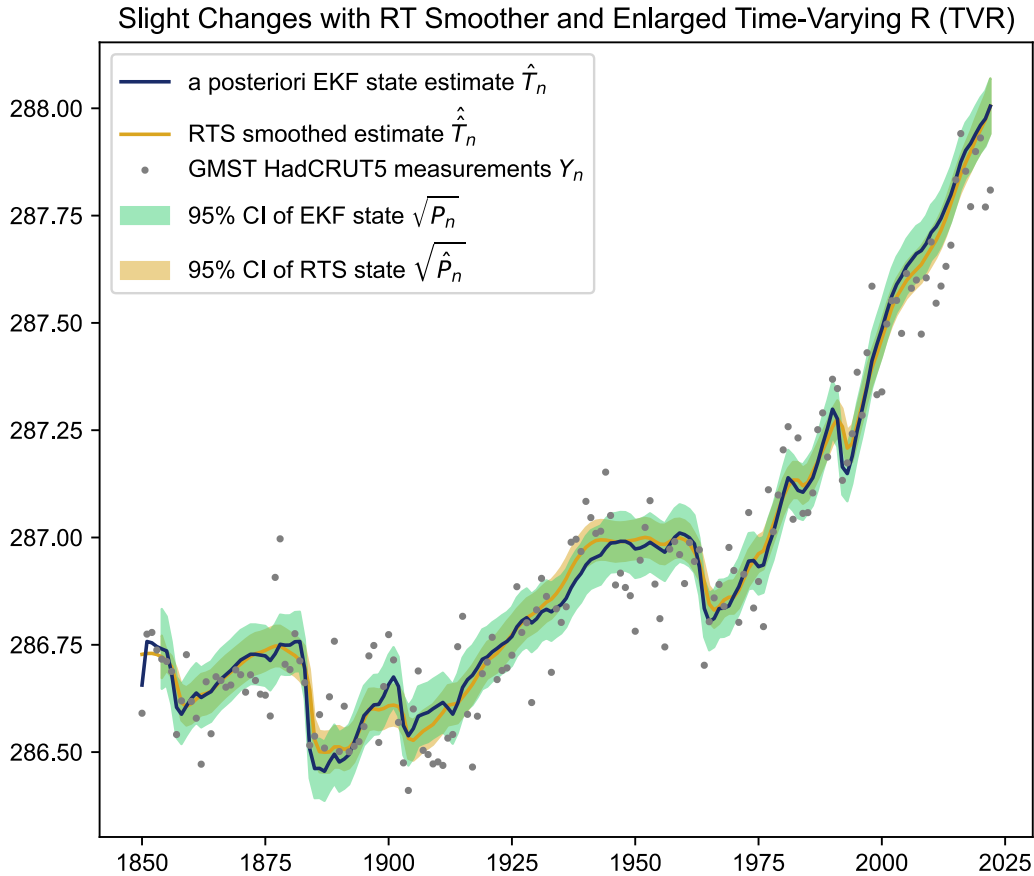
Section A: RTS Smoother

$$\widehat{K}_n = P_n \Phi_n (P_{n|n-1})^{-1} \quad \text{back-updated Kalman gain} \quad (\text{SA1})$$

$$\widehat{x}_n = \widehat{x}_n + \widehat{K}_n (\widehat{x}_n - \mathbf{F}(\widehat{x}_n; u_{n+1})) \quad \text{back-updated state estimate} \quad (\text{SA2})$$

$$\widehat{P}_n = P_n + \widehat{K}_n (\widehat{P}_{n+1} - P_{n|n-1}) \widehat{K}_n^T \quad \text{back-updated state covariance} \quad (\text{SA3})$$

This RTS has a theoretical advantage of blending abrupt changes in the model state over greater time periods, while also slightly reducing the state covariance. For instance, if the measurements suddenly and persistently diverged from the blind, forward EBM (unrelated to a known volcanic eruption), an EBM-Kalman Filter model state would only react as these measurements diverge, whereas an EBM-RTS would slightly foreshadow this jump because it can see future as well as past measurements. This occurred in 1943 and 1944: even though the EBM-KF estimated state is trending up, the EBM-RTS state moves cooler to reflect the colder GMST measurements from 1945-1950. Generally, the EBM-RTS just provides a second “nudge” toward measurements. Also, between 1850 and 1860 the intentionally overestimated initial state uncertainty P_0 of 1K is reduced through successive filtering steps in the EBM-Kalman Filter and bi-directional smoothing steps within the EBM-RTS. However, for the purposes of this paper, these distinctions make little difference, as is demonstrated in Supp. Fig. 1 below.

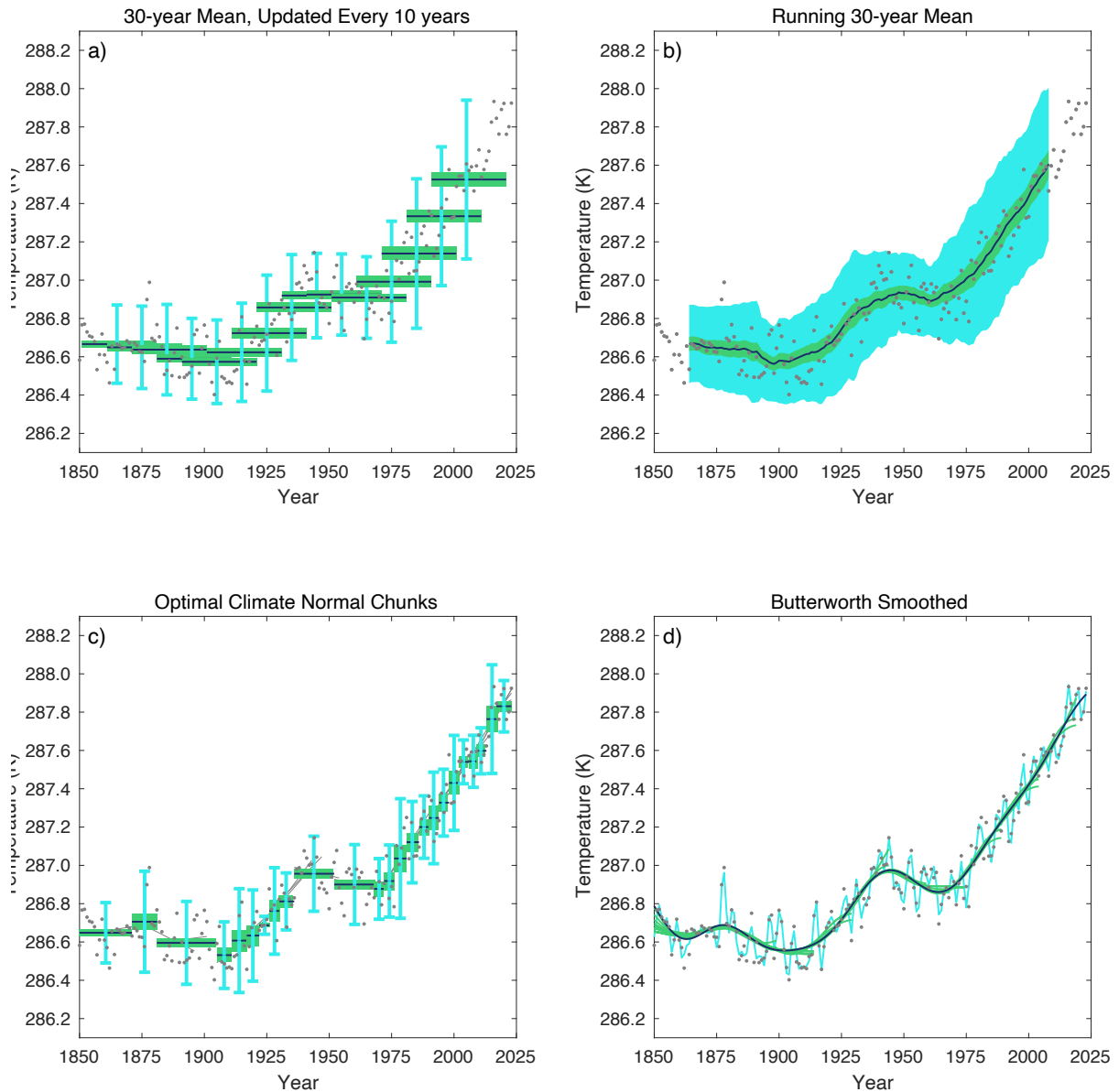


32

33 Supp. Fig. 1: Comparisons of the original EBM-Kalman Filtered climate state (navy blue line
 34 with green 1σ uncertainty window) with an EBM-RTS climate state (red line with red 1σ
 35 uncertainty window) and the effects of incorporating additional time-varying measurement
 36 uncertainty (green line with light blue 1σ uncertainty window). The addition of extra time-
 37 varying measurement uncertainty makes very little difference to the EBM-Kalman Filtered
 38 climate state, except from 1905-1930 when it lessens the deflection of repeated cooler GMST
 39 temperature measurements. In contrast, the EBM-RTS climate state doubly takes these
 40 annual temperature measurements into account, so it has a greater cooling deflection in this
 41 period, and many years that are warmer than the EBM-Kalman Filtered climate state after
 42 1980, although even these differences are slight - at most 0.1K during years of volcanic
 43 activity. However, there is greater certainty in the state estimate with the EBM-RTS: \hat{P}_n
 44 shrinks relative to P_n (see Supp. Fig. 10) by factors of 2.25 and 2.84 for the GMST and
 45 OHCA components respectively. The off-diagonal heat-transfer uncertainty component of \hat{P}_n
 46 is negative and 29 times smaller than those of P_n .

47

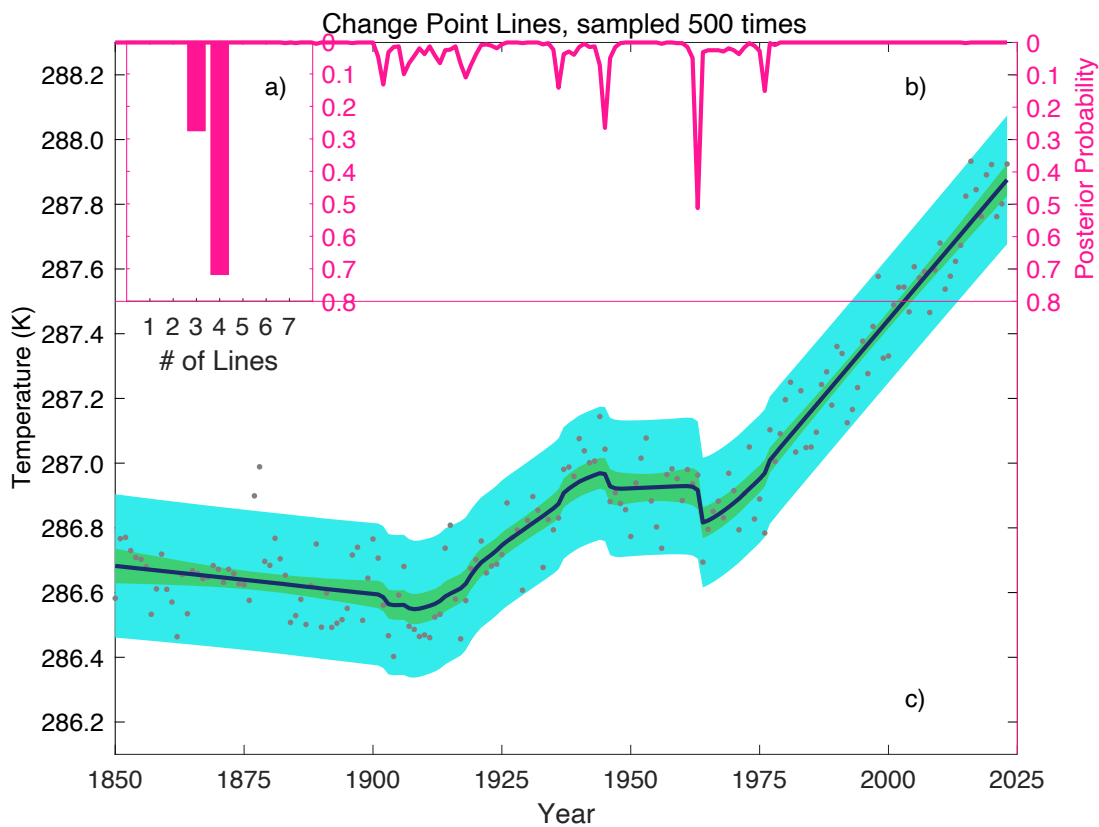
48



49

50 Supp. Fig. 2: Comparison of Prior Methods for Filtering or Smoothing the Climate as applied
 51 to the HadCRUT5 temperature dataset. (Morice, Kennedy et al. 2021) All metrics analogous
 52 to standard deviation are plotted at the 2σ level in light blue, and all metrics analogous to the
 53 standard error are plotted at the 1σ level in light green. a) The 30-year climate normals,
 54 updated every 10 years as per the World Meteorological Association in 1935. b) A running
 55 30-year average. c) Adaptive periods of multiyear averages, known as the optimal climate
 56 normal (OCN). (Livezey, Vinnikov et al. 2007). Chunks became smaller as the rate of climate
 57 change increased in recent decades. d) The Butterworth Smoother applied to this temperature
 58 dataset. (Mann 2008) For the “standard error” highly smoothed lines, the lowpass adaptive,
 59 lowpass mean padded, and lowpass methods were applied to chunks of the timeseries data
 60 ranging from 50 to 170 years in increments of 15 years with a cutoff frequency of $1/30$ years.
 61 The black “best” line a lowpass adaptive curve extended to 2021. The blue “standard
 62 deviation” line is a lowpass mean padded filter with a cutoff frequency of $1/5$ years.

63



65

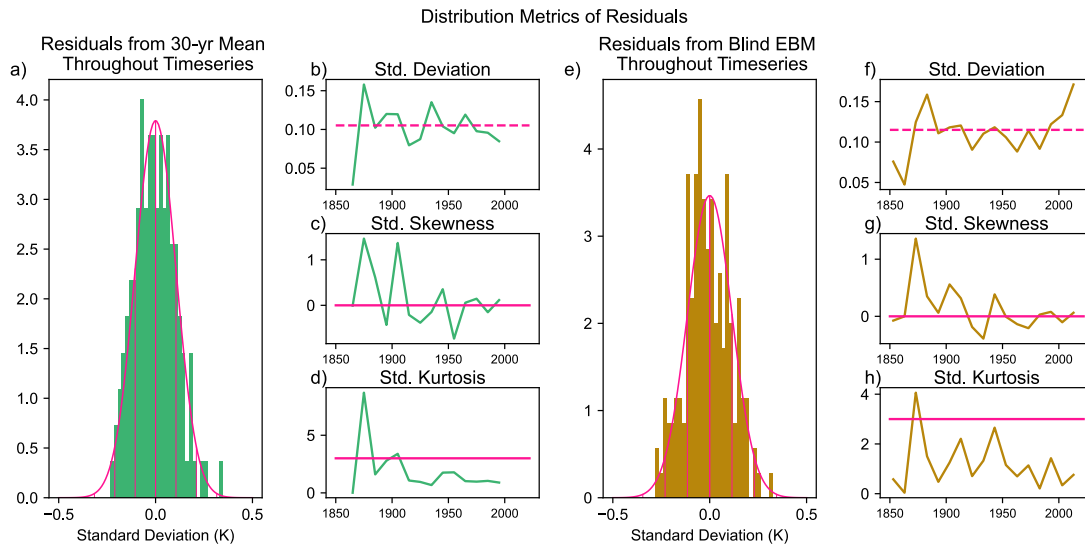
66 Supp. Fig. 3: Utilization of Bayesian Change Point on the HadCRUT5 data. (Ruggieri and
 67 Antonellis 2016) a) There are likely 4 trendlines with 72% of the posterior probability, and
 68 the remaining posterior probability on 3 trendlines. b) The posterior probability plot of where
 69 trendlines are most likely to occur: 51.2% of all samplings have a change point occur in 1963,
 70 and 26.4% of samplings have a change point occur in 1945. c) The posterior distribution of
 71 the trendlines in GMST, again with blue shading to indicate 2σ confidence interval of the
 72 data and green shading to indicate 2σ confidence interval of the mean trendline. These trend
 73 lines do not have to be continuous (note the dip at 1963), but over many samplings the
 74 average trend is smoothed.

75

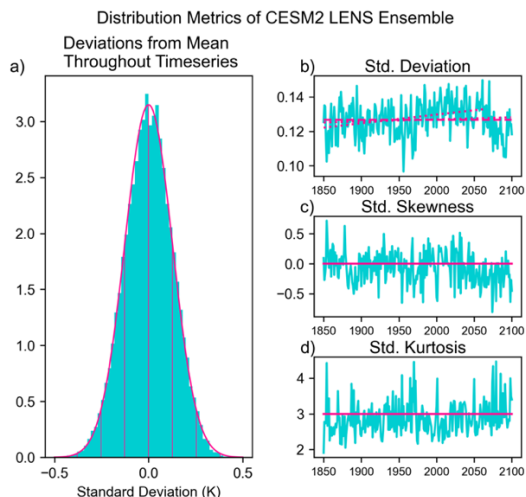
76

77

78

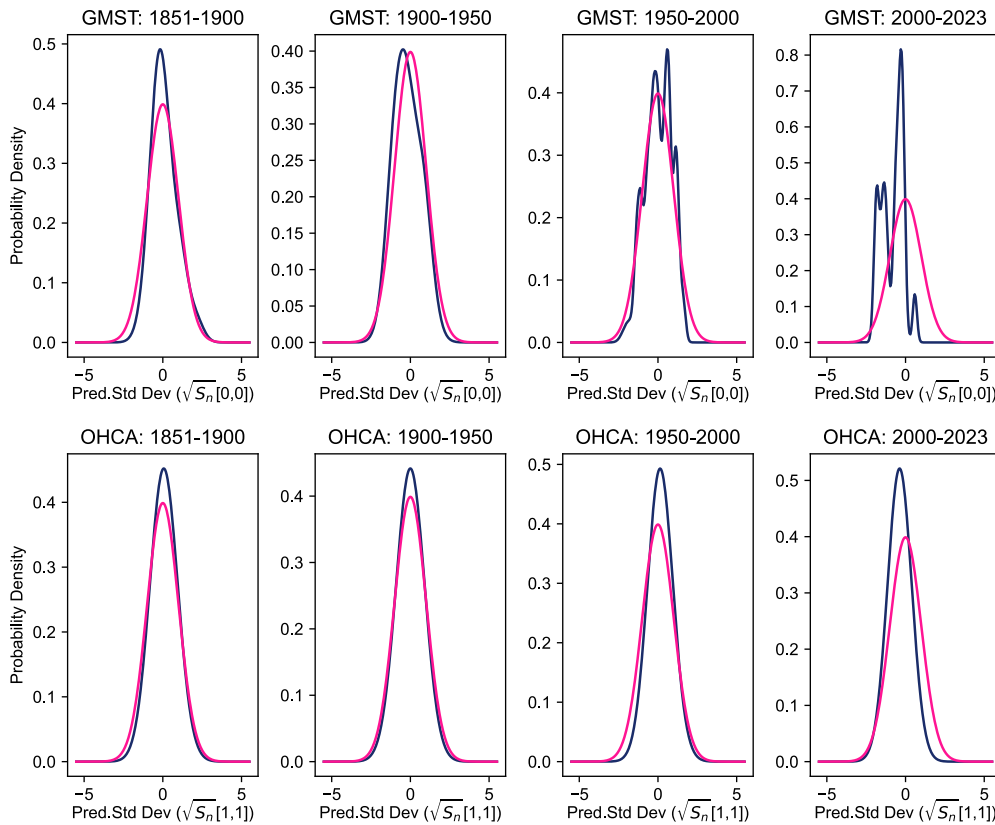


79
 80 Supp. Fig. 4: Left panels show statistical features of the residuals between the HadCRUT5
 81 measurements with respect to their 30-year running mean, which have a bias of -0.00339K .
 82 Pink lines in the histogram in (a) depict an ideal Gaussian distribution with standard deviation
 83 of 0.105K , and vertical lines drawn for each of these standard deviations. The dashed pink
 84 line (b) indicates the overall standard deviation. Solid pink lines for the skewness = 0.147 (c)
 85 and kurtosis = 1.904 (d) indicate the ideal values for a Gaussian distribution.
 86 Right panels show statistical features of the differences between the HadCRUT5
 87 measurements with respect to the blind EBM, which have a bias of -0.00104K . Pink lines in
 88 the histogram in (e) depict an ideal Gaussian distribution with standard deviation of 0.115K ,
 89 and vertical lines drawn for each of these standard deviations. The dashed pink line (f)
 90 indicates the overall standard deviation. Solid pink lines for the skewness = 0.123 (g) and
 91 kurtosis = 1.208 (h) indicate the ideal values for a Gaussian distribution.



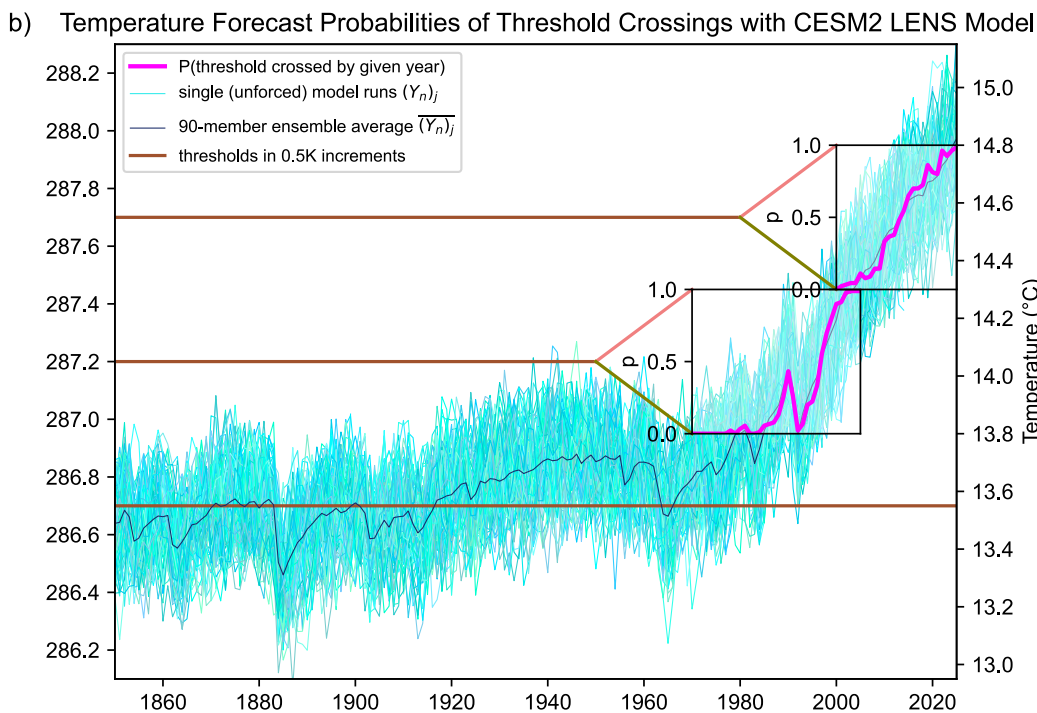
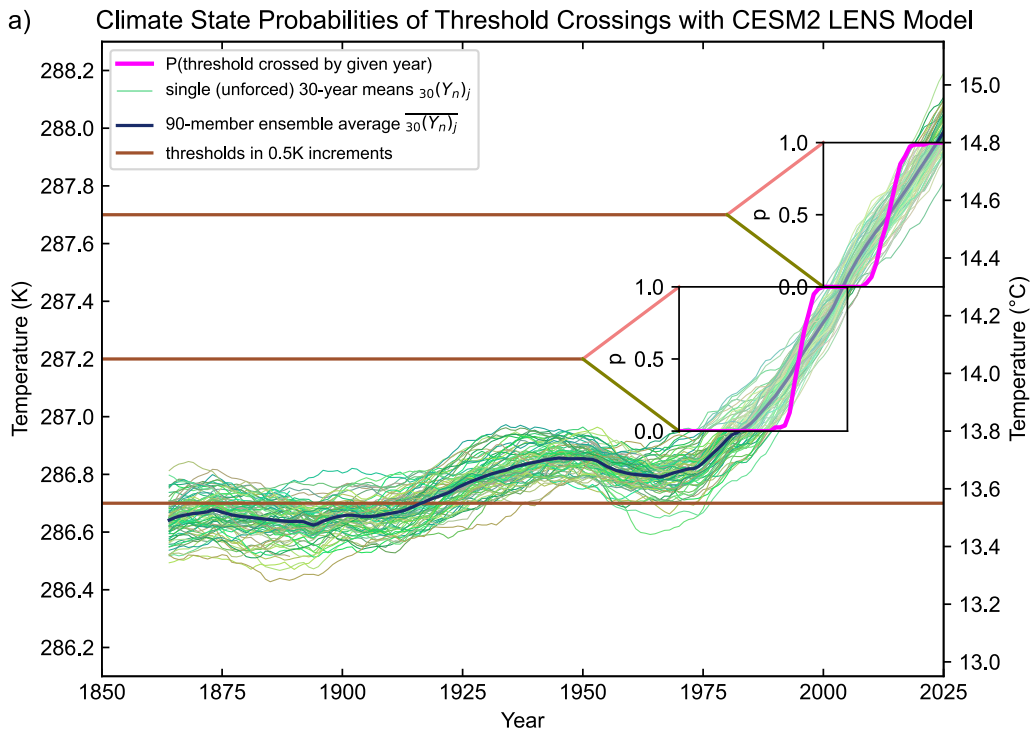
92
 93 Supp. Fig. 5: Statistical Features of the CESM2 Large Ensemble. (Rodgers, Lee et al. 2021).
 94 Pink lines in the histogram in (a) depict an ideal Gaussian distribution with standard deviation
 95 of 0.127K , and vertical lines drawn for each of these standard deviations. The observed trend
 96 (b) up until 2065 and overall in the standard deviation over time is plotted in a dotted pink
 97 line, while the dashed line indicates the overall standard deviation of 0.127K . Solid pink lines
 98 for the skewness = -0.067 (c) and kurtosis = 2.871 (d) indicate the ideal values for a Gaussian
 99 distribution.

EBM-KF Residuals Over Time



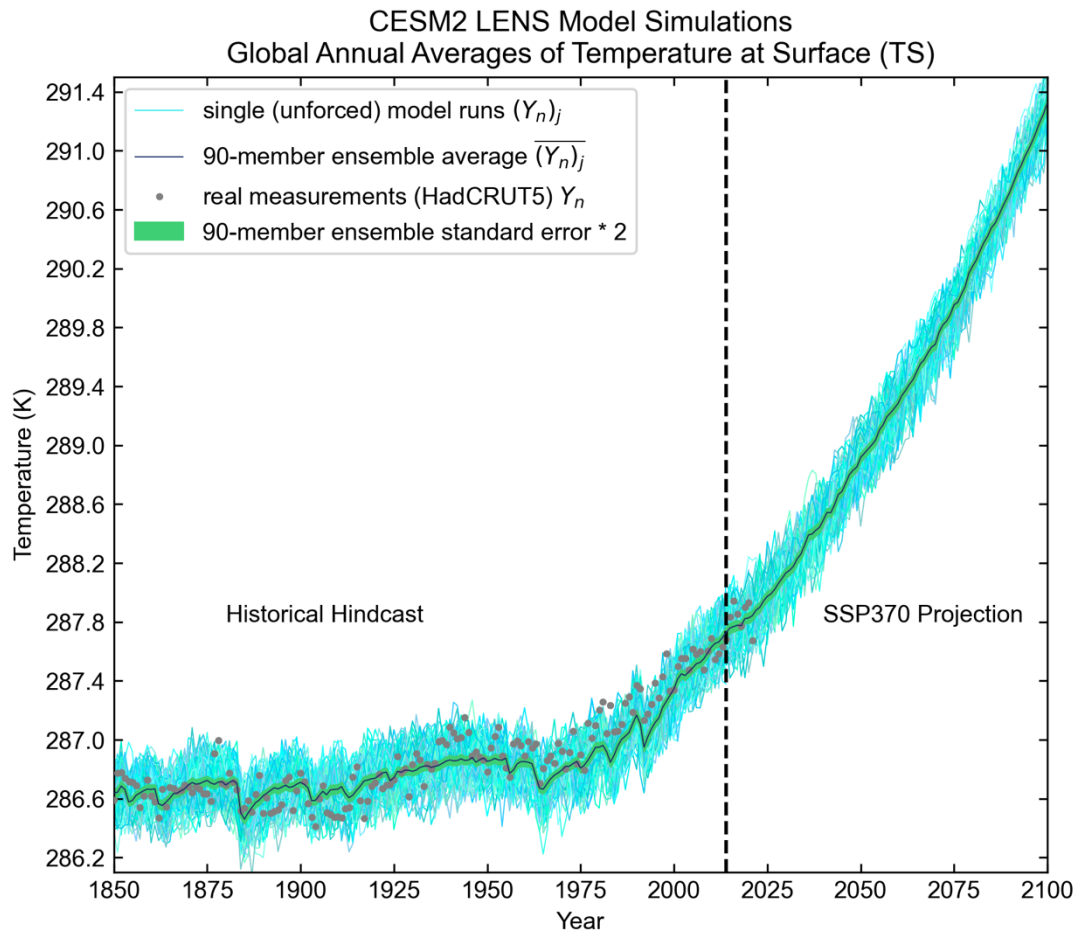
100

101 Supp. Fig. 6: Deviation between the projected climate state (pink) and empirical PDFs of the
 102 Gaussian mixture of measurements with associated uncertainty (purple), plotted relative to
 103 the ideal distribution given by the innovation covariance. Each column indicates a different
 104 time window of the EMB-KF model's run length. The top row displays the empirical PDFs of
 105 the GMST HadCRUT measurements relative to the model's estimate of GMST state, whereas
 106 the bottom row displays empirical PDFs of the OHCA Zanna 2019 measurements relative to
 107 the model's estimate of OHCA state. Note the initial period begins at 1851 (and the 1850
 108 measurement is excluded from main text Fig. 3 and 4) because this has comparison involves
 109 P_0 , which was intentionally over-estimated (resulting in relatively too-narrow measurement
 110 kernel). Also note that the last period is less than half the time of the others, so the GMST
 111 empirical distribution is much more choppy.

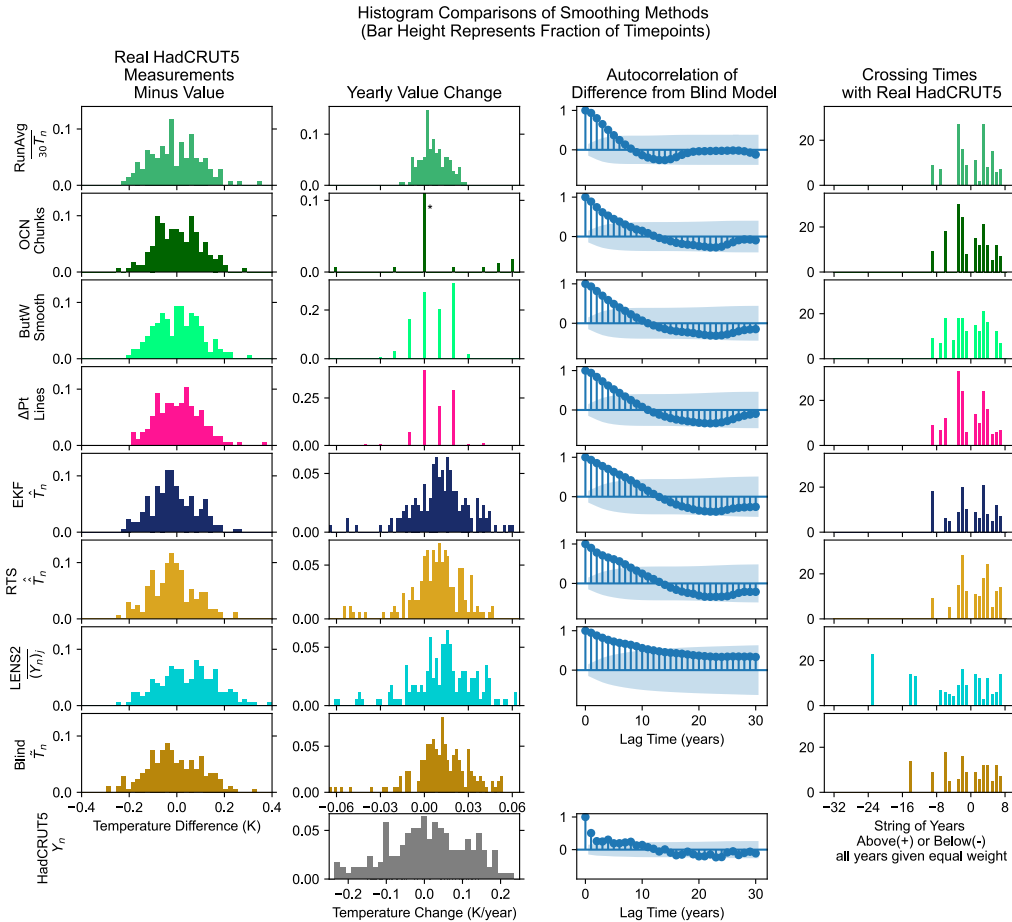


112

113 Supp. Fig. 7: As in Fig. 5 within the main text, except the threshold crossing calculations are
 114 performed on the LENS2 ensemble. a) climate state policy thresholds: the 21-year running
 115 means of individual simulations in light green lines, the two inset boxes indicate threshold
 116 crossing probability, given by the fraction of these light green lines that have crossed the
 117 indicated threshold. b) temperature forecast policy thresholds, showing a cloud of the
 118 possible next-year measurements in light blue from the simulations, and again the two inset
 119 boxes indicate the fraction of these light blue lines that have crossed the threshold.

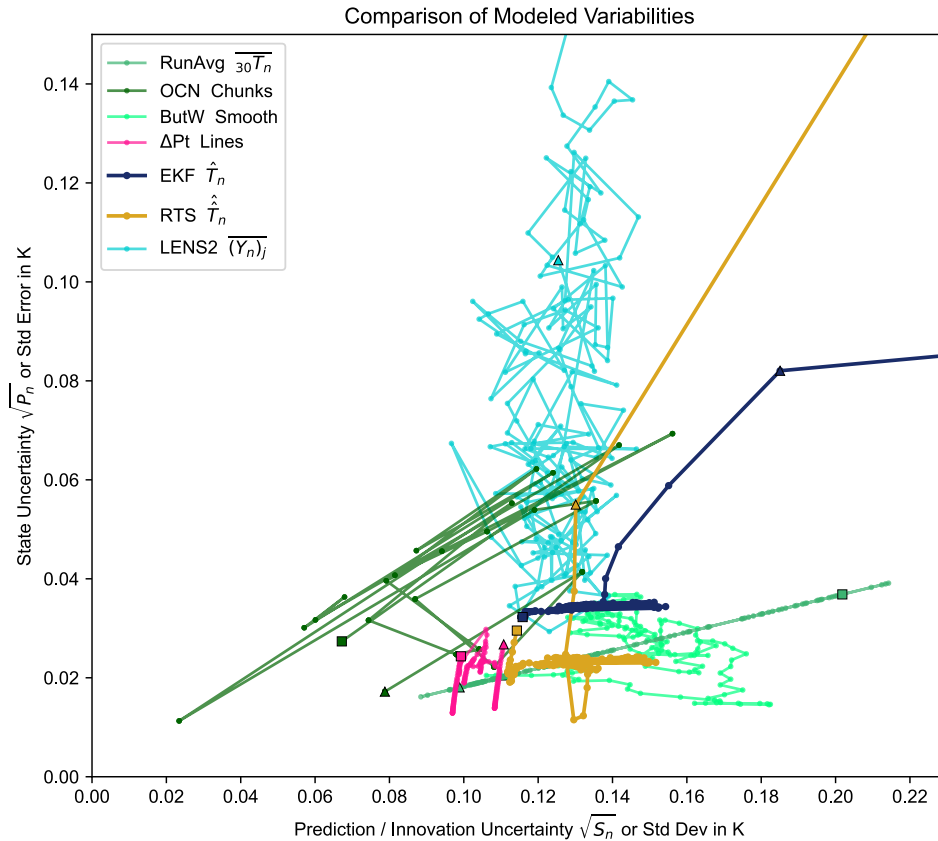


121
 122 Supp. Fig. 8: Comparison of the CESM2 Large Ensemble (LENS2) GSAT (Rodgers, Lee et
 123 al. 2021) with HadCRUT5 GMST measurements. The various shades of thin light blue and
 124 turquoise lines represent each individual simulation $(Y_n)_j$ of the 90-member ensemble. The
 125 ensemble mean is plotted in a navy-blue line, and the ensemble mean standard error is plotted
 126 around this line in green. This standard error is twice the standard deviation divided by the
 127 square root of the number of ensemble members at that moment and shows the 2σ uncertainty
 128 in the yearly simulated climate is roughly 0.026K. The ensemble mean has $R^2 = 0.83$ relative
 129 to the HadCRUT5 measurements, lower than for the blind EBM ($R^2=0.88$). The dashed
 130 vertical line represents when LENS transitions from historical to future forcing (SSP3-7.0).
 131
 132



133

134 Supp. Fig. 9: Histogram comparisons of several aspects of many of the smoothing methods
 135 for generating a climate timeseries. The far-left column represents the absolute differences
 136 between the HadCRUT5 measurements and all the other models. All look similar in this
 137 respect. The center-left column shows the annual changes in the temperatures reported by
 138 each model. In this respect, the real HadCRUT5 measurements are the most spread out,
 139 because the stochastic change each year is large, whereas the in most years the OCN Chunks
 140 do not change. The center-right column shows an autocorrelation plot, which demonstrates
 141 that every other model aside from HadCRUT5 (and to a lesser extent the running average) are
 142 autocorrelated with the blind energy-balance model to similar degrees. The far-right column
 143 shows how many continuous years are spent above or below HadCRUT5: both the LENS2
 144 ensemble average and the blind energy-balance model had >20 year spans for which they
 145 were colder than the “real” HadCRUT5 data, illustrating the benefit of data assimilation.
 146



147
 148 Supp. Fig. 10: Comparisons of the state and prediction (or equivalent) uncertainties of the
 149 smoothing methods for generating a climate timeseries. The x-axis represents the state
 150 uncertainty (colored light green in all other figures), and the y-axis represents the prediction
 151 uncertainty (colored light blue and doubled in all other figures). As these quantities change
 152 over time, all points in these smoothing timeseries are connected with colored lines, with the
 153 triangle Δ representing the value of these quantities in 1850 or the first point that they entered
 154 the frame limits of this graph, and the square \square representing the value of these quantities in
 155 2021 or the last point that they were within the frame limits. For instance, the running
 156 average draws a straight line because standard deviation and standard error are linearly
 157 correlated by a factor of $1/\sqrt{30}$, and latter points have larger quantities for each variability due
 158 to the changing climate. The Butterworth Smoother traces a curve roughly in this region, with
 159 both the standard deviations and standard errors being twice the 15-year running average of
 160 the maximum of the absolute value of differences between colored and black curve. The RTS
 161 and EKF methods rapidly converge to a state uncertainty of $\sim 0.110\text{K}$ and $\sim 0.03\text{K}$. The
 162 Change Point Regression variance also fluctuate in this region, although this methods'
 163 standard error twice drops to 0.014K . Both the OCN and the LENS2 climates have standard
 164 errors that are above the other methods at most times. For LENS2, the standard deviation
 165 within the CESM2 ensemble generally remains between 0.11K and 0.14K , whereas the state
 166 uncertainty is taken to be the standard deviation of the 20 ensembles comprising [CMIP6](#) in
 167 October 2021. (Meehl, Moss et al. 2014) These metrics have nothing to do with Figure 10 in
 168 the main text. Within CMIP6, the 20 ensembles are most in agreement in 1939, when the
 169 state uncertainty dipped down to only 0.029K between ensemble means, but this uncertainty
 170 was much greater at earlier and later time points, reaching 0.183K by 2014.
 171
 172

173 **Section C: Justification that the EKF is sufficient, will not diverge**

174 The issue of nonlinearity arises not in the computation of $\hat{x}_{n|n-1}=F(\hat{x}_{n-1})$ but rather the
 175 covariance distribution P_n of points (infinitesimal probability masses) neighboring \hat{x}_{n-1} , which
 176 are assumed to scale linearly around this transformation to maintain a normal distribution.
 177 The OHCA part of the model can be ignored since it is purely linear. Nonlinear distortion
 178 may pile more probability density onto a state other than the transformed original projection
 179 $F(\hat{x}_{n-1})$, necessitating a new computation of $\hat{x}_{n|n-1}$ as the mean of this distorted PDF. Thus, for
 180 an arbitrary point that is z standard deviations away from \hat{x}_{n-1} , the remainder error R_1
 181 (Lagrange mean-value form) induced in a single cycle is:

$$182 \quad F(\hat{x}_{n-1}+z\sqrt{P_n};u_n) - F(\hat{x}_{n-1}) - \frac{\partial F(x;u_n)}{\partial x} z\sqrt{P_n} =$$

$$183 \quad R_1(\hat{x}_{n-1}+z\sqrt{P_n}) = \frac{\partial^2 F(\xi_L;u_n)}{\partial \xi_L^2} \frac{(z\sqrt{P_n})^2}{2} \quad \text{for } \xi_L \in [\hat{x}_{n-1}-|z|\sqrt{P_n}, \hat{x}_{n-1}+|z|\sqrt{P_n}] \quad (\text{SC1})$$

$$184 \quad = \left(\frac{0.441}{\text{AOD}_n+9.73} (0.00159) - (0.00005546) (1 - 0.0655 \log_{10}([\text{CO}_2]_n)) 1.385 (\xi_L)^{0.385} \right) \frac{z^2 P_n}{2}$$

$$185 \quad (\text{SC2})$$

$$186 \quad -0.5(10^{-5}) z^2 P_n < R_1(\hat{x}_{n-1}+z\sqrt{P_n}) < 0.5 (10^{-5}) z^2 P_n \quad (\text{SC3})$$

$$187 \quad |R_1(\hat{x}_{n-1}+z\sqrt{P_n})| < 10^{-5} z^2 0.5 (0.032)^2 < |z| * 5 * 10^{-9} \quad (\text{SC4})$$

188 This means that all probability masses that are within $|z| < 20$ standard deviations will have an
 189 one-step error of $< 0.000002\text{K}$. Even if the error accumulates in the same direction in each
 190 cycle of the EKF, over the 173 year timeseries, the error will be within 0.0004K compared to
 191 a particle method such as the Unscented Kalman Filter. (Julier and Uhlmann 1997; Wan and
 192 Van Der Merwe 2000)

193